

International Workshop on
High Energy Density Physics
Beijing, April 19 - 20, 2010

Two lectures by J. Meyer-ter-Vehn :

1. On fast ignition: Selfsimilar isochoric implosions

reflects my scientific work on Inertial Fusion (1980 -1995)

2. On laser-driven ultra-thin foils as relativistic mirrors

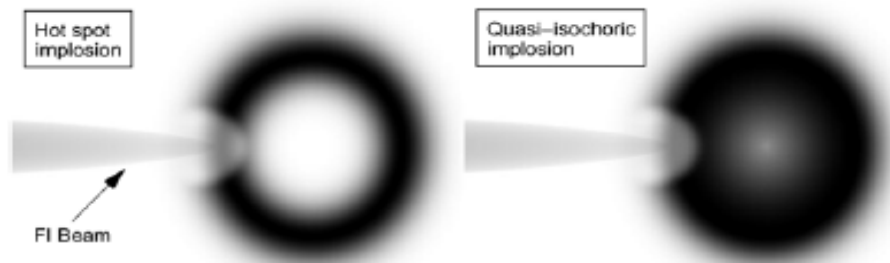
reflects my work on Relativistic Laser Plasma (1995 - 2010)

On fast ignition: Selfsimilar isochoric implosions

J. Meyer-ter-Vehn, MPQ Garching

hot spot
ignition

fast
ignition

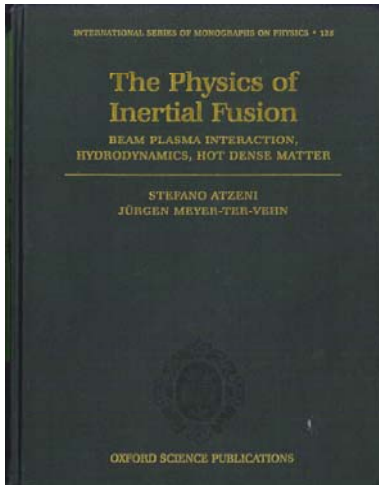


A self-similar isochoric implosion for fast ignition
M.S. Clarke and M.Tabak, Nucl.Fus. 47, 1147 (2007)

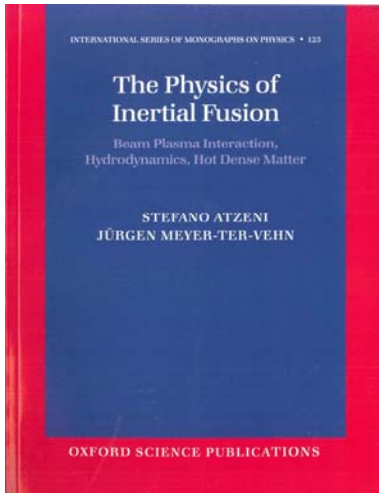


The Book

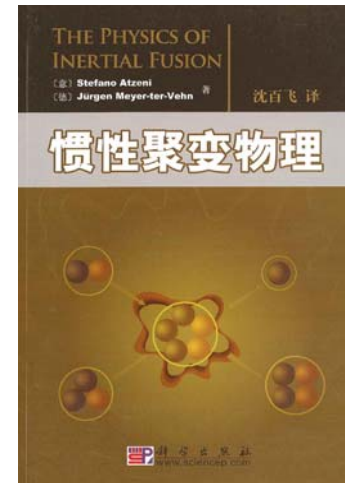
The Physics of Inertial Fusion
S. Atzeni & J. Meyer-ter-Vehn



Oxford University
Press 2004



Paperback 2009



Chinese Edition
2009
translated by

Baifei Shen (沈百飞)



Equations of 1D ideal gas dynamics and symmetries

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^{n-1}} \frac{\partial}{\partial r} (r^{n-1} \rho u) = 0$$

Parameters : $n = 1, 2, 3$ $\gamma = 5/3$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial r} + \frac{1}{\rho} \frac{\partial}{\partial r} \left(\frac{\rho c^2}{\gamma} \right) = 0$$

The generators operate on objects

$$\{t, r, u, c, \rho\}$$

$$\left(\frac{\partial}{\partial t} + u \frac{\partial}{\partial r} \right) \left(\frac{\rho c^2}{\rho^\gamma} \right) = 0$$

Example for stretching :

$$\exp(\varepsilon r \partial_r) F(r) = F(e^\varepsilon r)$$

time translation $G_1 = \partial_t$

time stretching $G_2 = t \partial_t - u \partial_u - c \partial_c$

space stretching $G_3 = r \partial_r + u \partial_u + c \partial_c - n \rho \partial_\rho$

mass stretching $G_4 = \rho \partial_\rho$

} scaling
symmetries

projection $G_5 = t^2 \partial_t + r t \partial_r + (r - u t) \partial_u + c t \partial_c - n \rho t \partial_\rho$ (for $n = 2/(\rho - 1)$)
satisfied for $n=3, \gamma=5/3$

Similarity solutions of gas dynamics

Similarity solutions (Guderley 1942)

velocity $u(r,t) = (\alpha r/t)U(\xi)$

sound velocity $c(r,t) = (\alpha r/t)C(\xi)$

density $\rho(r,t) = \rho_0 (r/r_0)^\kappa G(\xi)$

Similarity coordinate:

$$\xi = \frac{r/r_0}{(t/t_0)^\alpha}$$

Two free parameters:

$$\alpha, \kappa,$$

are invariant under general scaling generator

$$G = t\hat{\partial}_t + \alpha r\hat{\partial}_r + (\alpha - 1)u\hat{\partial}_u + (\alpha - 1)c\hat{\partial}_c + \alpha\kappa\rho\hat{\partial}_\rho$$

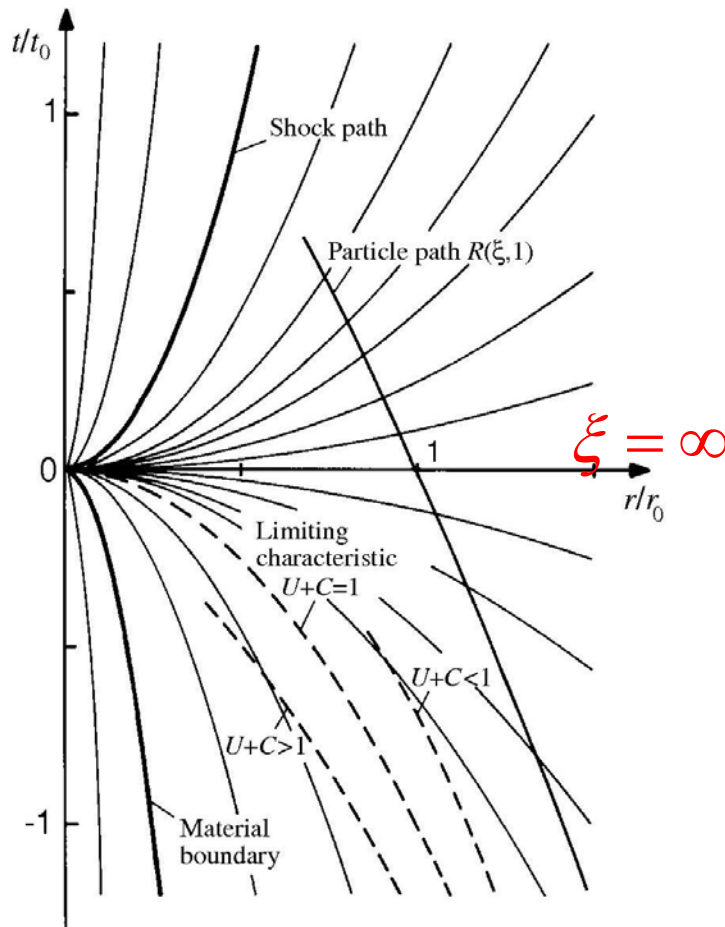
It allows to reduce the gas dynamics equations to

$$\frac{dU}{dC} = \frac{\Delta_1(U, C)}{\Delta_2(U, C)}$$

$$\frac{d \ln \xi}{dC} = \frac{\Delta(U, C)}{\Delta_2(U, C)}$$

$$G = G(U, C, \xi)$$

The similarity coordinate $\xi = r/|t|^\alpha$



Profiles at $t = 0$ ($\xi = \infty$)

$$u(r) = u_0 r^\lambda \quad (\lambda = 1 - 1/\alpha)$$

$$c(r) = c_0 r^\lambda$$

$$\rho(r) = \rho_0 r^\kappa$$

$$p(r) = p_0 r^{\kappa - 2\lambda}$$

Entropy

$$p / \rho^\gamma = A_0 r^\varepsilon \quad (\varepsilon = \kappa(\gamma - 1) + 2\lambda)$$

isentropic $\varepsilon = 0$

Mach number

$$M = u_0 / c_0$$

Guderley solution

G. Guderley,
Luftfahrtforschung 19, 302 (1942)

$$\frac{dU}{dC} = \frac{\Delta_1(U, C)}{\Delta_2(U, C)}$$

solve to get $U(C)$

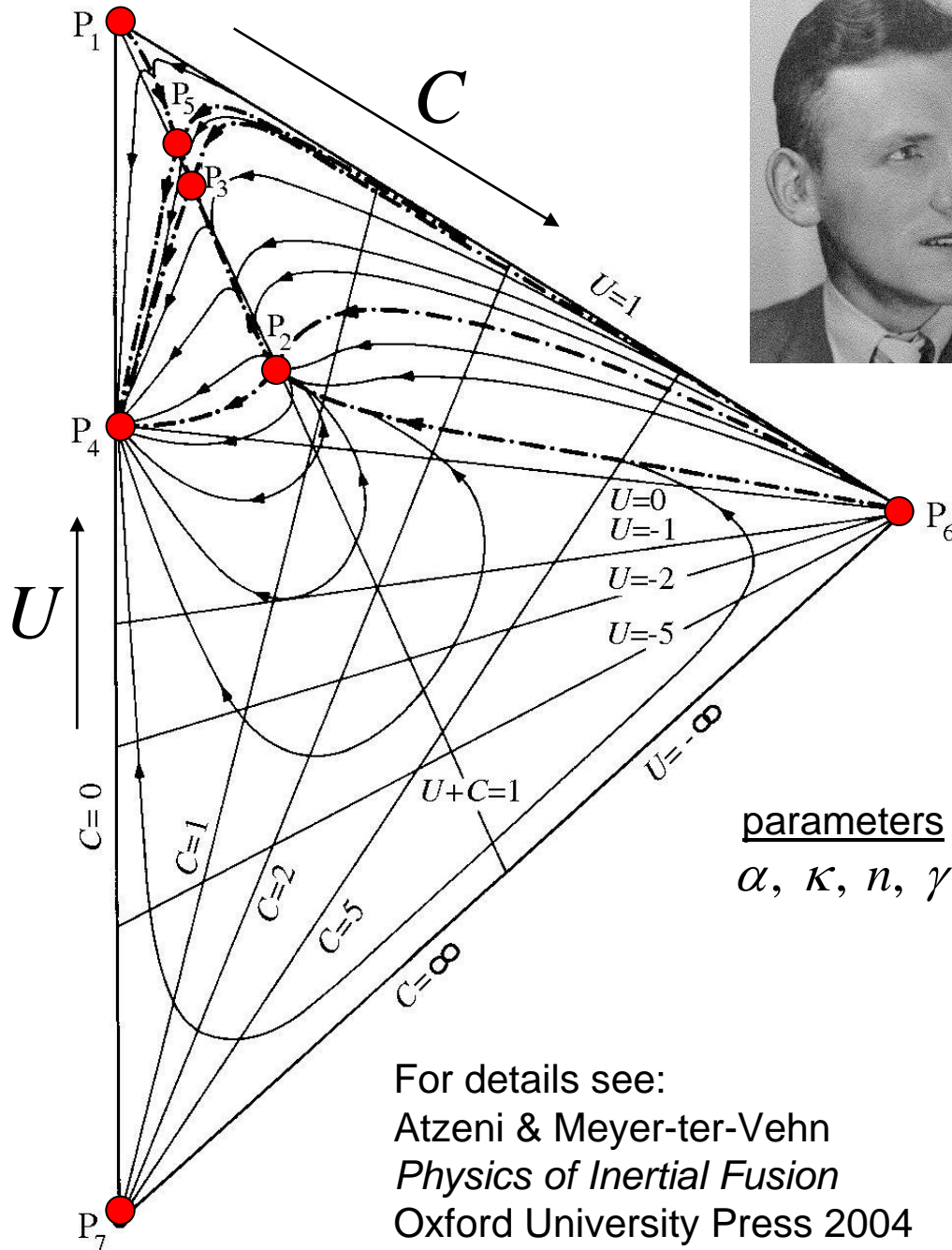
$$\frac{d \ln \xi}{dC} = \frac{\Delta(U, C)}{\Delta_2(U, C)}$$

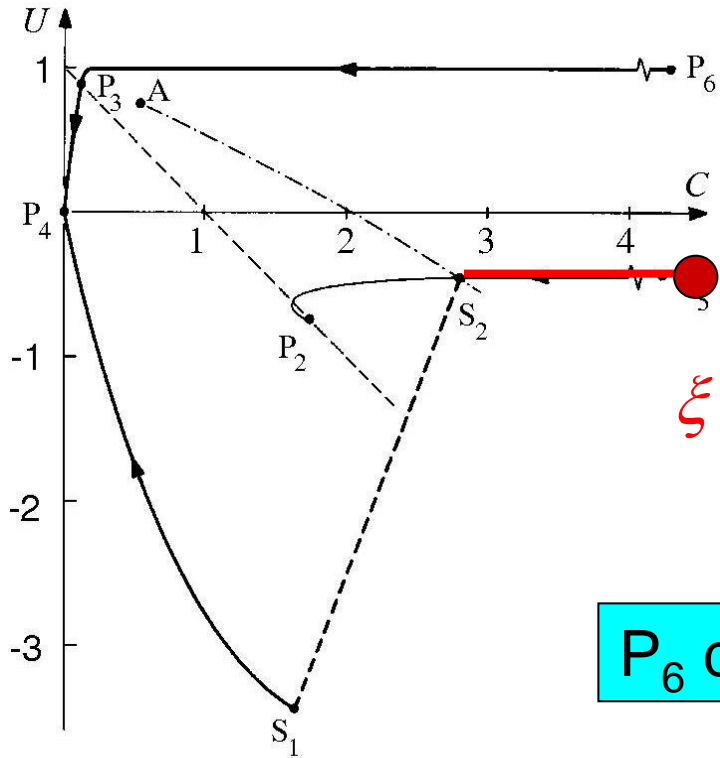
integrate to get $\xi(C)$

$$C = C(\xi)$$

$$U = U(\xi)$$

$$G = G(U, C, \xi)$$





$$P_6: \quad C_6 = \infty, \quad U_6 = \frac{2(1-1/\alpha) - \kappa}{n\gamma}$$

$$\xi = 0$$

P_6 describes gas state in center

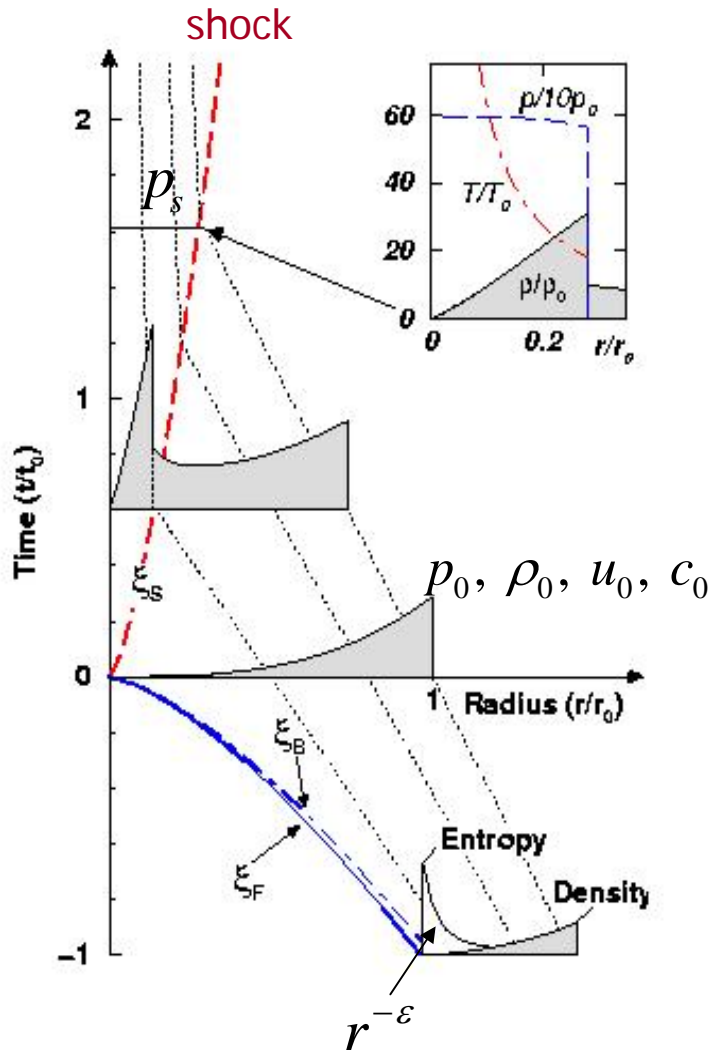
$$u(r, t) = \frac{2(1-1/\alpha) - \kappa}{n\gamma} \alpha \frac{r}{t}$$

$$\rho(r, t) \sim r^{n\varepsilon/v} t^{\alpha(\kappa - n\varepsilon/v)}$$

$$T(r, t) \sim c^2 \sim r^{-n\varepsilon/v} t^{\alpha(-2\lambda + n\varepsilon/v)}$$

$$p(r, t) \sim r^0 t^{\alpha(\kappa - 2\lambda)} \quad (v = n\gamma + \kappa - 2\lambda)$$

Mach number and entropy determine implosion



nearly uniform pressure
behind shock p_s

at void closure ($t=0$)

$$M = u_0 / c_0 = \text{const}$$

key result

$$p_s / p_0 \cong 3.6 M^3$$

$$\alpha = 0.7$$

$$\kappa = 3$$

$$\varepsilon = \kappa(\gamma - 1) + 2(1/\alpha - 1) = 2.85$$

Stagnation pressure scaling

Meyer-ter-Vehn, Schalk, Z. Naturforschung 37a, 955 (1982)

Kemp, Meyer-ter-Vehn, Atzeni, PRL 86, 3336 (2001).

$$\frac{p_s}{p_0} \sim M^{\frac{2(n+1)}{\gamma+1}} \sim M^3 \quad \left\{ \begin{array}{l} \text{for} \\ n=3 \\ \gamma=5/3 \end{array} \right.$$

Mach number

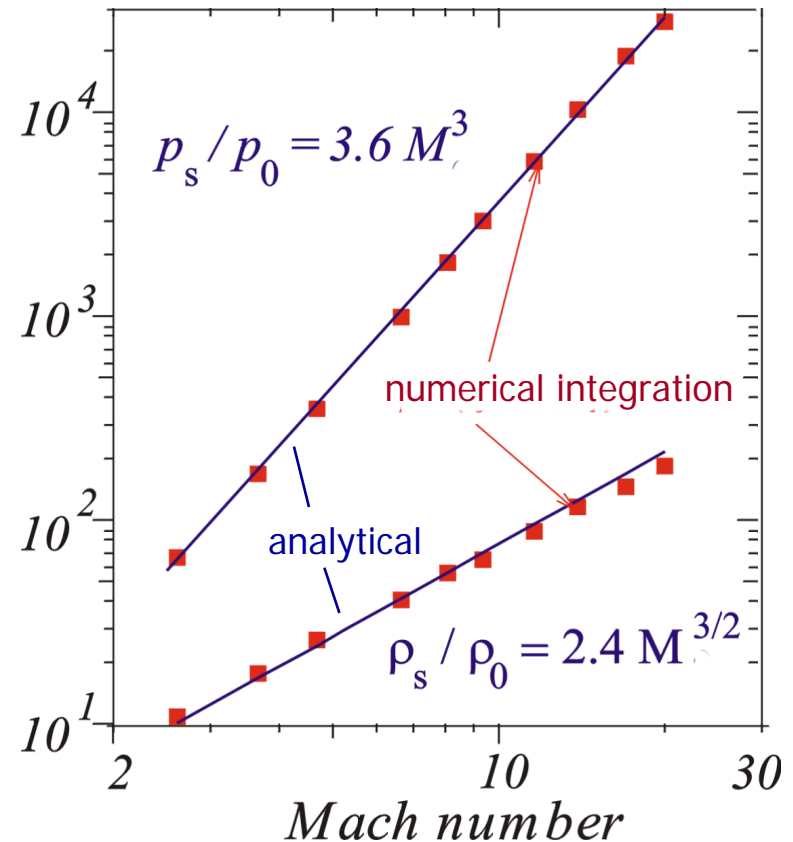
$$M = u_{imp} / c_0$$

$$c_0 \sim (p_0 / \rho_0)^{1/2}$$

$$p_0 \sim \alpha_{if} \rho_0^{5/3}$$

Stagnation pressure

$$p_s \sim u_{imp}^3 \alpha_{if}^{-9/10} p_a^{2/5}$$



Ignition energy scaling

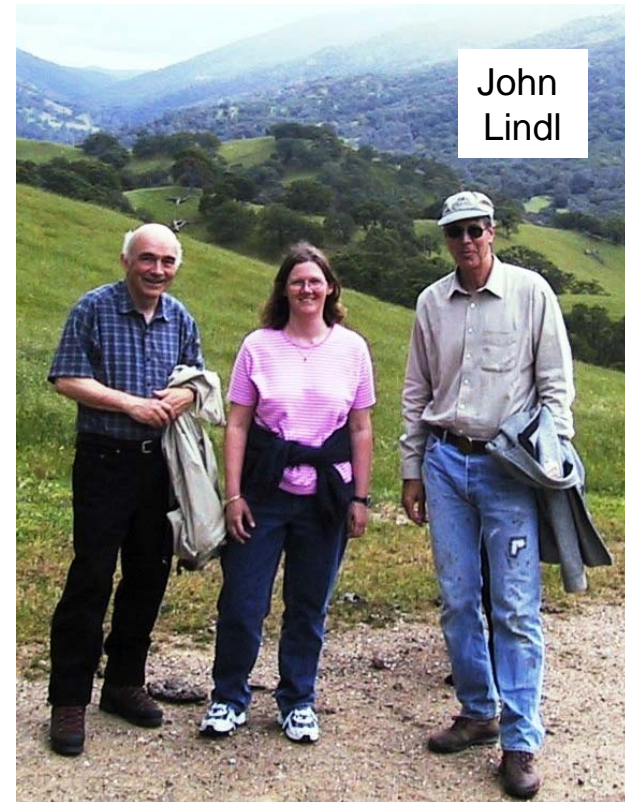
$$E_f^{ign} \sim \frac{1}{\rho_s^2} \sim \frac{1}{\rho_a^2 M^6} \sim \alpha_{if}^{1.8} u_{imp}^{-6} \rho_a^{-0.8}$$

Atzeni, Meyer-ter-Vehn, Nucl.Fus. 41, 465 (2001)

$$E_f^{ign} \sim \alpha_{if}^{1.88 \pm 0.05} u_{imp}^{-5.89 \pm 0.12} \rho_a^{-0.77 \pm 0.12}$$

scaling extracted from large set of simulations

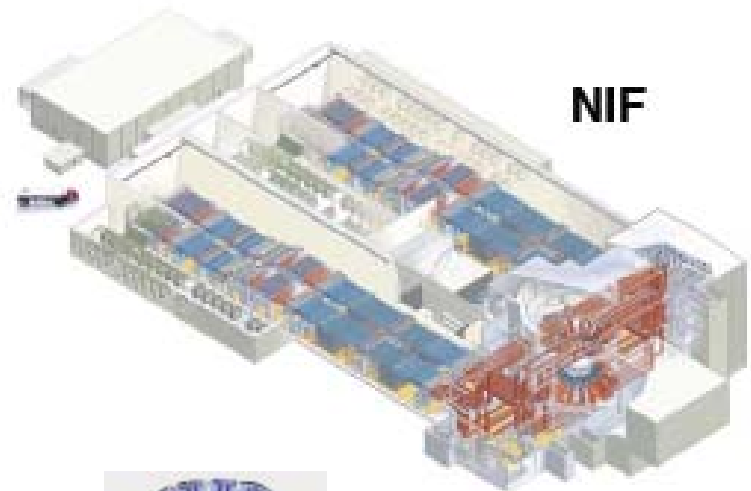
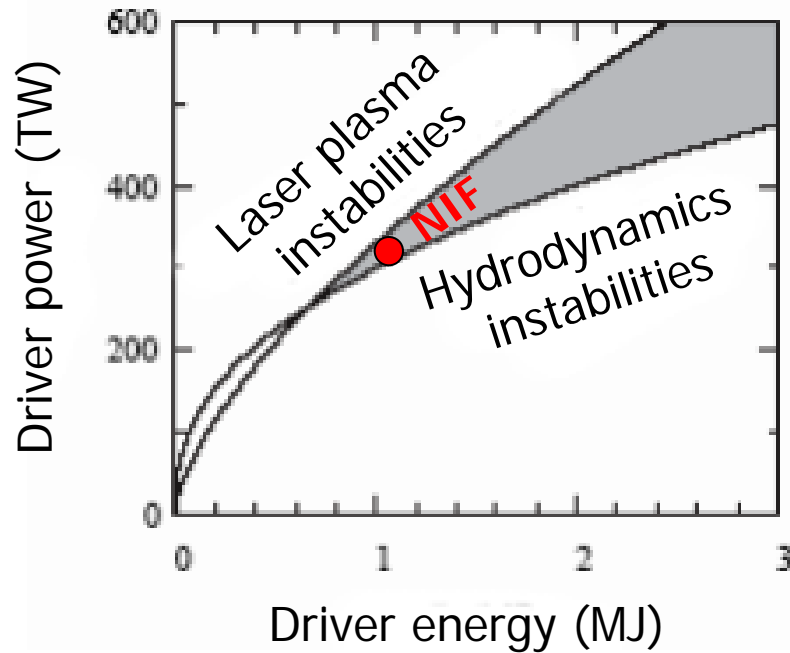
Herrmann, Tabak, Lindl, Nucl.Fus. 41, 99 (2001)



The ICF Ignition Window

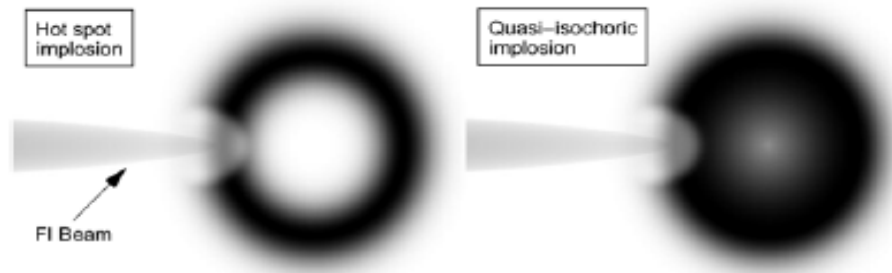
J. Lindl et al. *Physics of Plasmas* 11, 339 (2004)

S. Atzeni & J. Meyer-ter-Vehn *Physics of Inertial Fusion* Oxford University Press (2004)

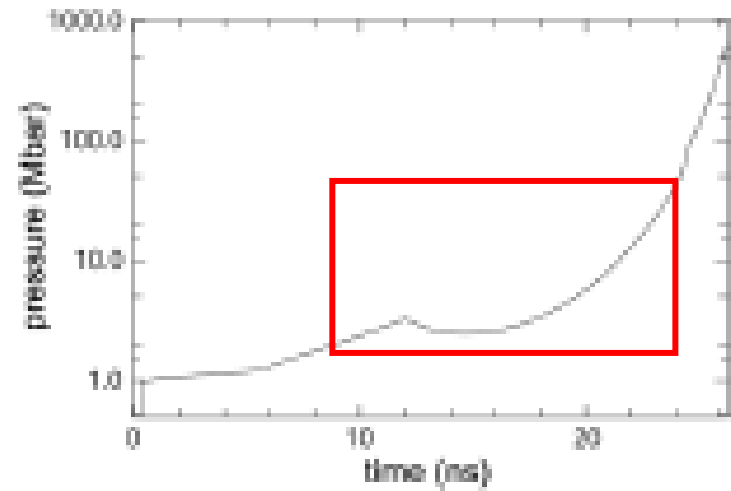
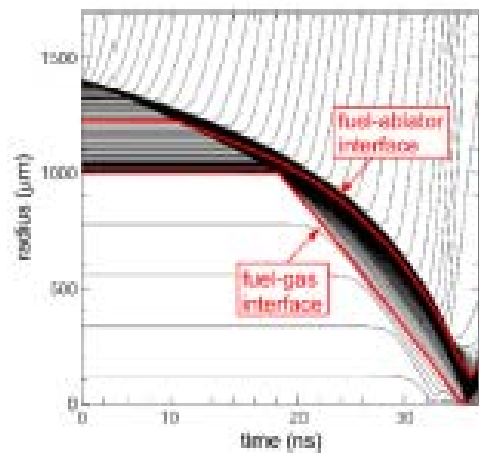


A self-similar isochoric implosion for fast ignition

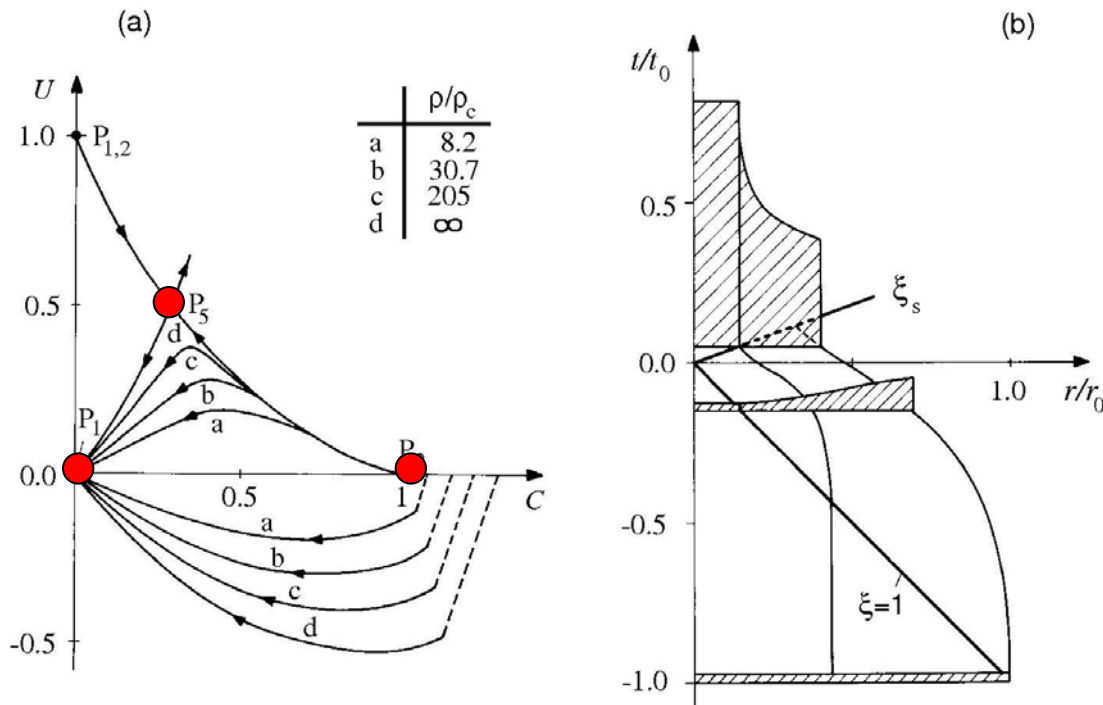
D.S. Clark and M. Tabak



Look for solutions with $\epsilon = 0$



The ideal isochoric implosion :
 Compressing a uniform gas at rest
 into a uniform gas at rest of arbitrary density



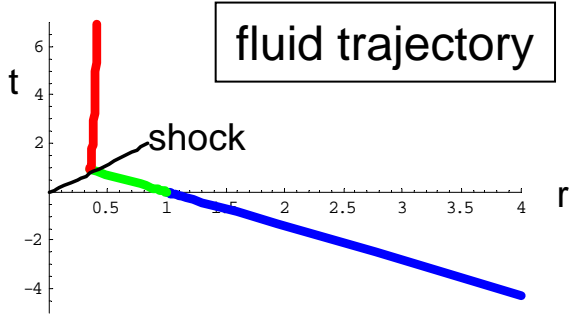
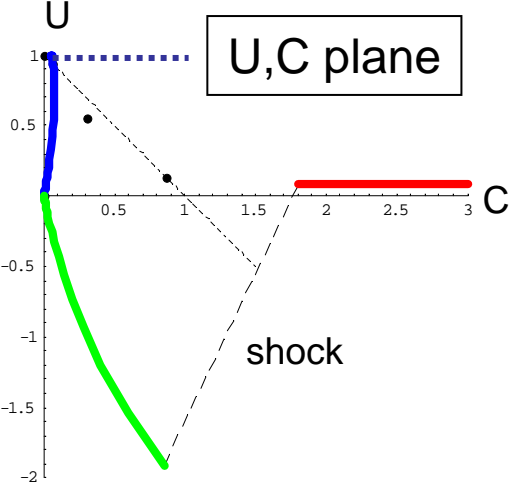
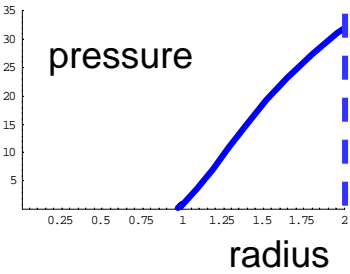
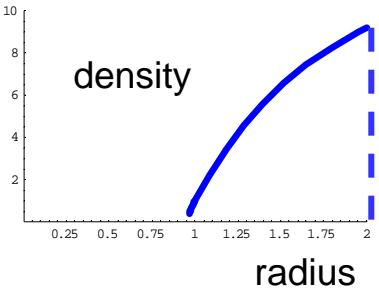
Disadvantage:

Low Mach number,
 return shock weak,

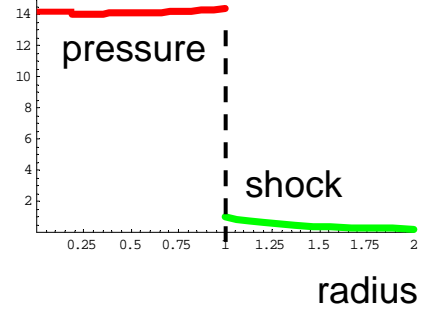
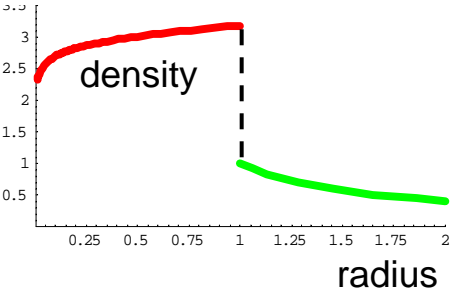
Almost full pressure
 of igniting fuel
 (>100 Gbar)
 has to be balanced
 by external piston

High Mach number, nearly isochoric solution: $\alpha = 0.9, \varepsilon = 0.1$

Imploding shell



Stagnating shell



almost constant
implosion velocity

Projection symmetry holds for $n = 2/(\gamma - 1)$

Book (Chapter 6.6)

$$G_p = t^2 \partial_t + rt \partial_r + (r - ut) \partial_u - ct \partial_c - n \rho t \partial_\rho$$

allows to generate new solution from known solution

$$\exp(\varepsilon G_p) \{r, t, u, c, \rho\} \{r^{(0)}, t^{(0)}, u^{(0)}, c^{(0)}, \rho^{(0)}\} = \{r^{(\varepsilon)}, t^{(\varepsilon)}, u^{(\varepsilon)}, c^{(\varepsilon)}, \rho^{(\varepsilon)}\}$$

old solution

$$\xi = \frac{r}{|t|^\alpha}$$

$$u(r, t) = \alpha \frac{r}{t} U(\xi)$$

$$c(r, t) = \alpha \frac{r}{t} C(\xi)$$

$$\rho(r, t) = \rho_0 r^\kappa G(\xi)$$

projected new solution

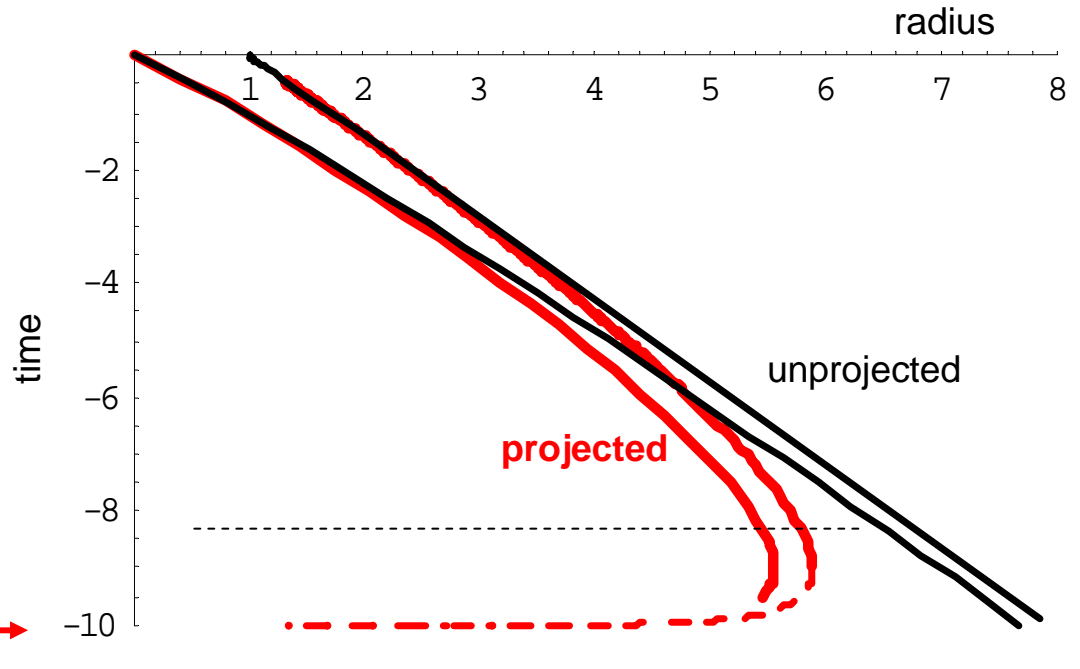
$$\xi = \frac{r}{|t|^\alpha (1 + t/t_0)^{1-\alpha}}$$

$$u(r, t) = \frac{r}{t(1 + t/t_0)} [\alpha U(\xi) + t/t_0]$$

$$c(r, t) = \frac{r}{t(1 + t/t_0)} \alpha C(\xi)$$

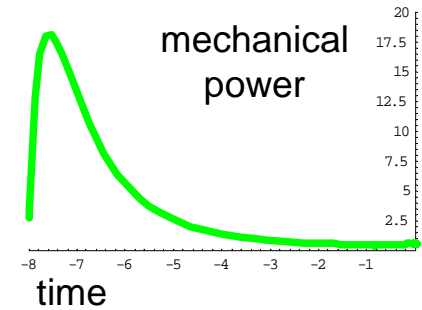
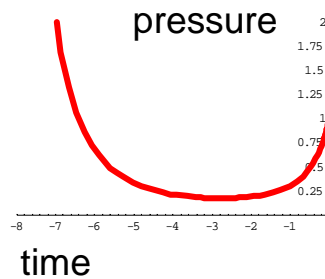
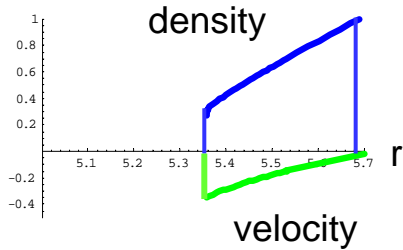
$$\rho(r, t) = \rho_0 \frac{r^\kappa}{t(1 + t/t_0)^{n+\kappa}} G(\xi)$$

Projected solution



profiles at $t = -8$

applied to outer trajectory:



Conclusions

Similarity solutions provide deep insight into the gas dynamics of ICF implosions

They also serve as a guide for isochoric implosions in the context of fast ignition